
Micromagnetic 3D simulation of the pinning field in high temperature $\text{Sm}(\text{Co,Fe,Cu,Zr})_2$ magnets

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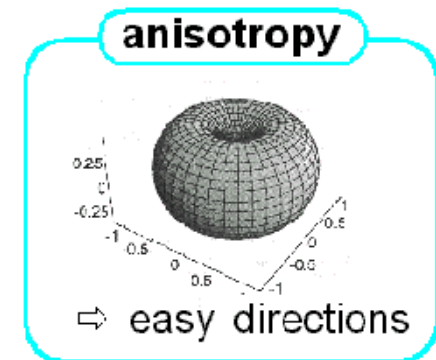
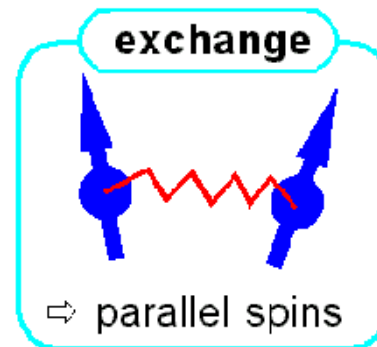
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Outline

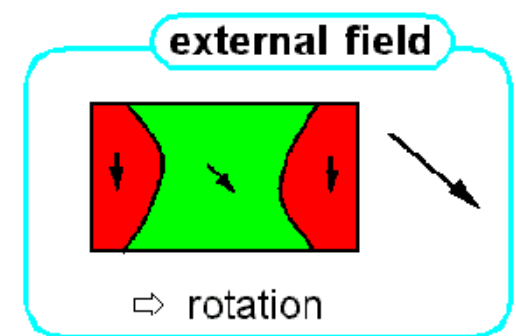
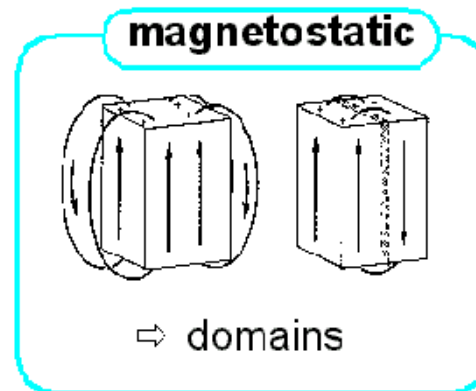
- **Pinning Controlled SmCo Magnets**
- **Micromagnetic Model**
- **Attractive Pinning**
- **Repulsive Pinning**
- **Variation of the Phase Thickness**
- **Summary**

Micromagnetics

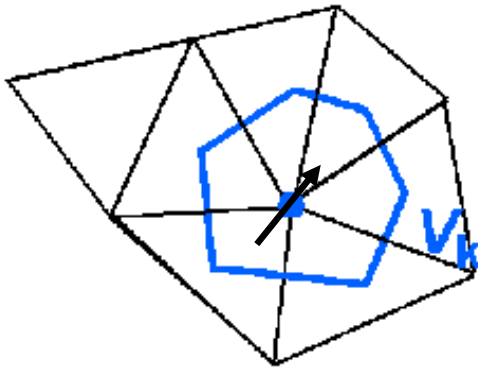
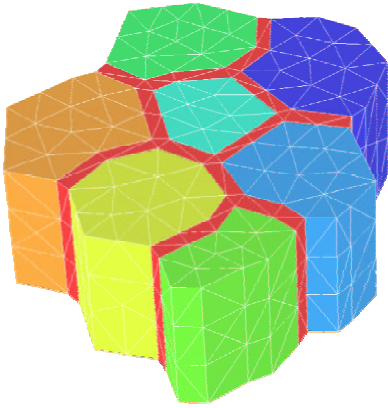
- Effective field H_{eff} :
 - exchange
 - anisotropy
 - magnetostatic
 - external field
- Find energy minimums by integration of the Gilbert equation of motion or direct energy minimization



$$\frac{\partial \mathbf{J}}{\partial t} = -|\gamma| \mathbf{J} \times \mathbf{H}_{\text{eff}} + \frac{\alpha}{J_s} \mathbf{J} \times \frac{\partial \mathbf{J}}{\partial t}$$



Finite Element Approach



- divide particles into finite elements
⇒ triangles, tetrahedrons
- expand \mathbf{J} with basis function φ

$$\vec{J}(\vec{x}) = \sum_{i=1}^{nodes} \vec{J}_i \varphi_i(\vec{x})$$

- energy as a function of $\mathbf{J}_1, \mathbf{J}_2 \dots \mathbf{J}_N$

$$E(\vec{J}_1, \vec{J}_2 \dots \vec{J}_N)$$

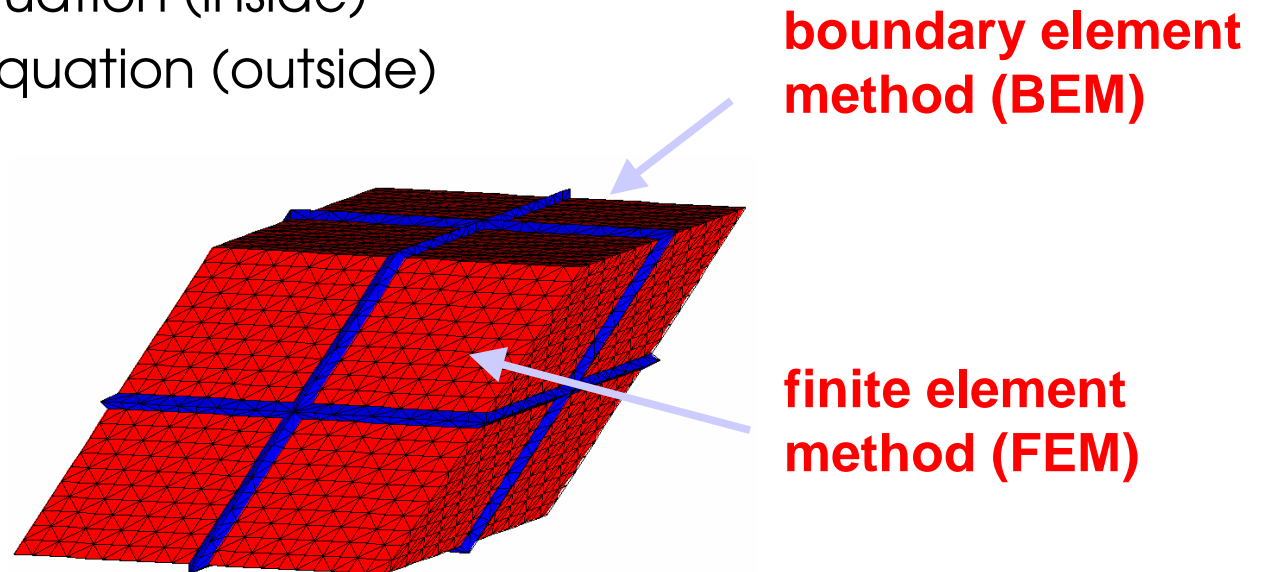
- effective field

$$\vec{H}_k = -\frac{1}{V_k} \frac{\partial E(\vec{J}_1, \vec{J}_2 \dots \vec{J}_N)}{\partial \vec{J}_k}$$

- ⇒ effective field on irregular grids
- ⇒ rigid magnetic moment
at the **nodes**

Magnetostatic Field Calculation

- ▶ magnetic scalar potential
 $\mathbf{H} = -\nabla U$
- ▶ solve Gilbert equation *simultaneously* with
 - ⇒ Poisson equation (inside)
 - ⇒ Laplace equation (outside)



BEM leads to a fully populated $N \times N$ matrix

- ⇒ N ... number of nodes at the surface
- ⇒ matrix compression using wavelets

Pinning Controlled SmCo Magnets

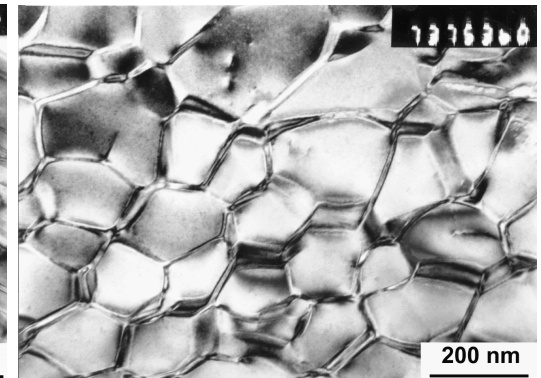
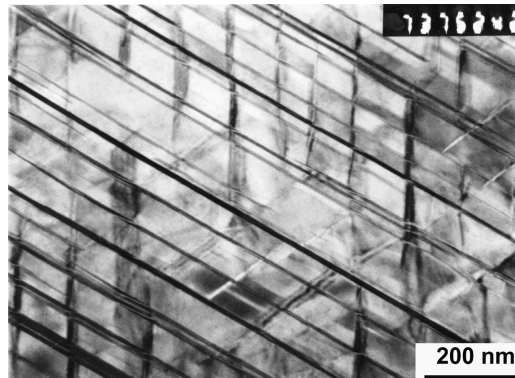
- Characterization of SmCo permanent magnets by transmission electron microscopy

Microstructure

- Composition
- Heat treatment
- Additives

influence

- Precipitation structure
- Lamella phase
- Cell size

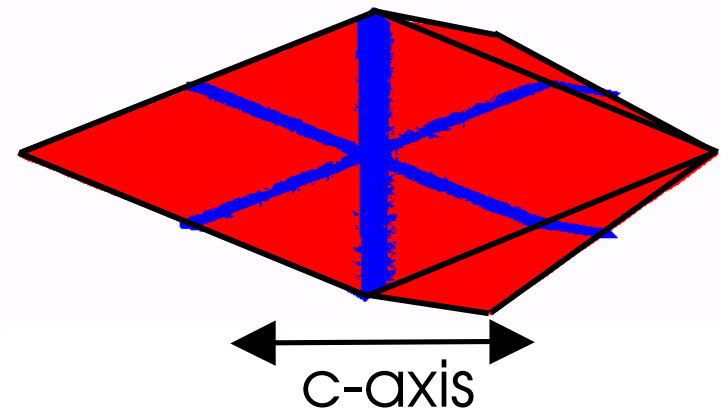
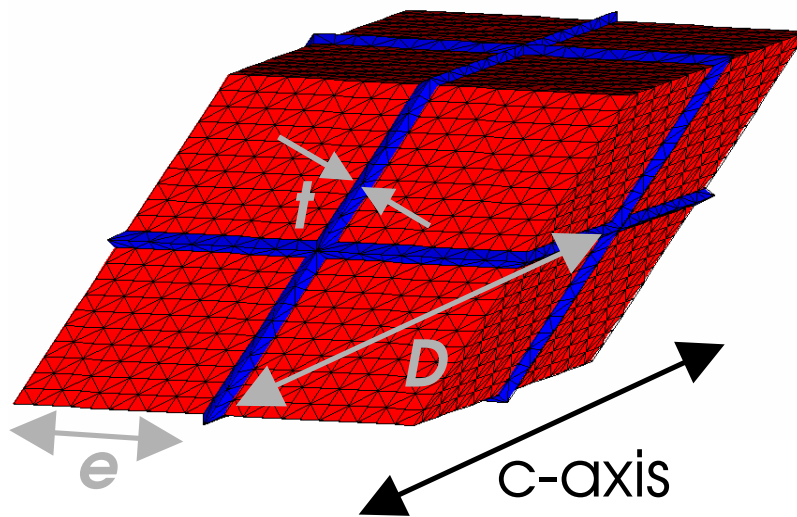


- Lorentz image of two magnetic domains

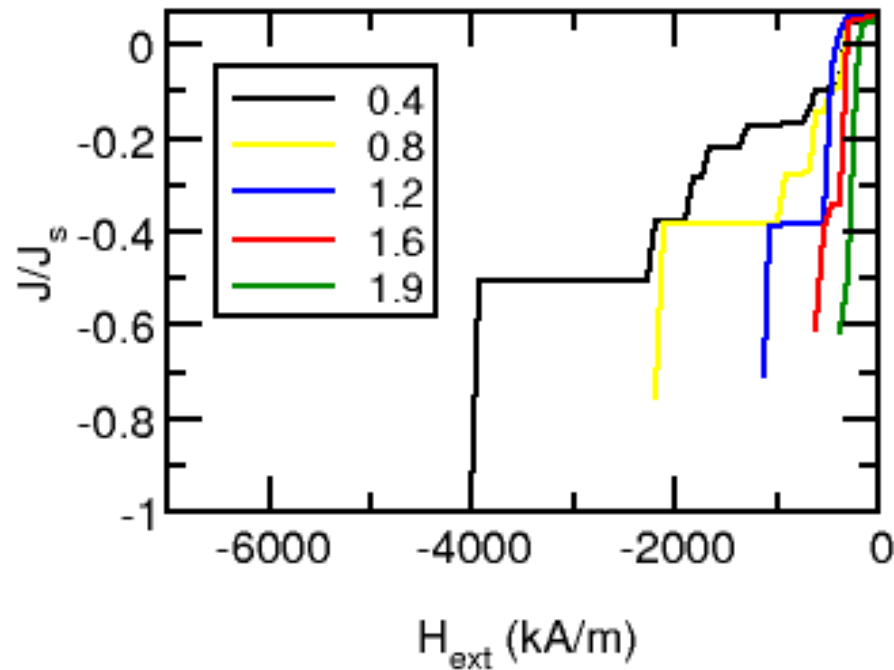


Micromagnetic Model

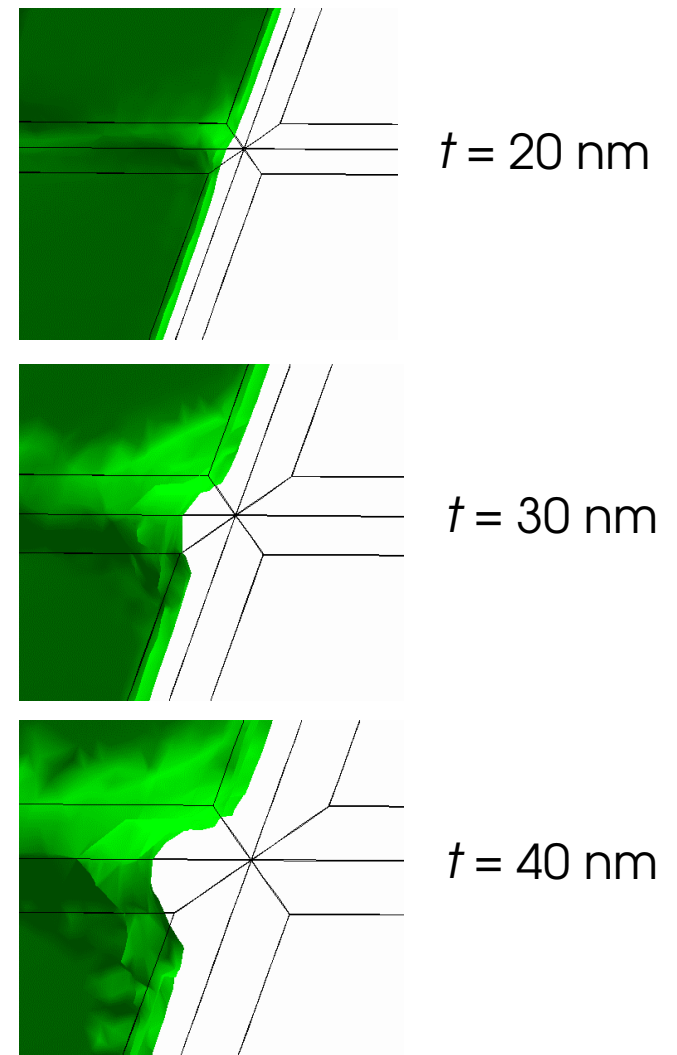
- Finite element mesh:
15833 nodes
84749 tetrahedral elements
7056 surface elements
- Resolution of the mesh:
 $e/10 = D/25$
for $D = 125$ nm: 5 nm
 $\delta(\text{Sm}_2\text{Co}_{17}) = 5$ nm



Attractive Pinning



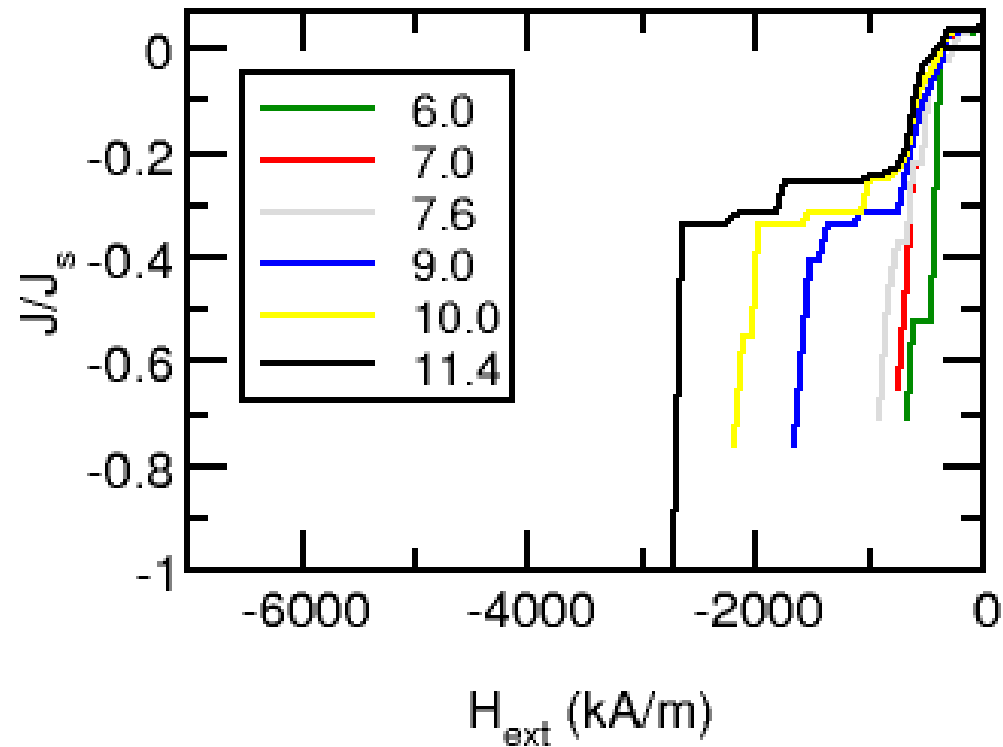
Domain wall bending



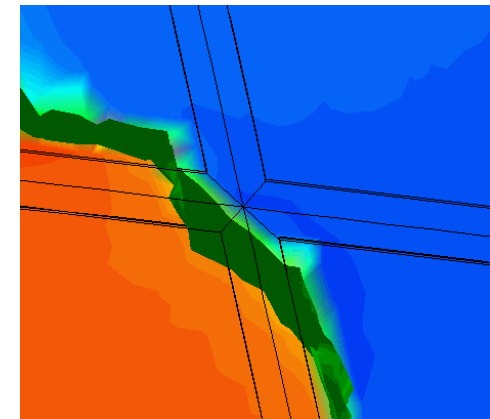
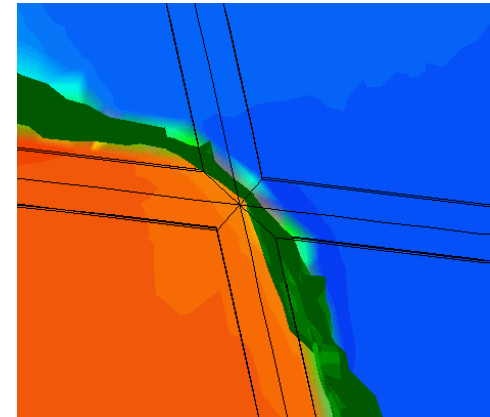
- Cells ($\text{Sm}_2\text{Co}_{17}$):
Polarization: $J_s = 1.32$ T
Anisotropy: $K_1 = 5$ MJ/m³
Exchange: $A = 14$ pJ/m
- Intercellular phase:
Polarization: $J_s = 0.8$ T
Anisotropy: $K_1 = 1.2$ MJ/m³
Exchange: $A = 14$ pJ/m

Repulsive Pinning

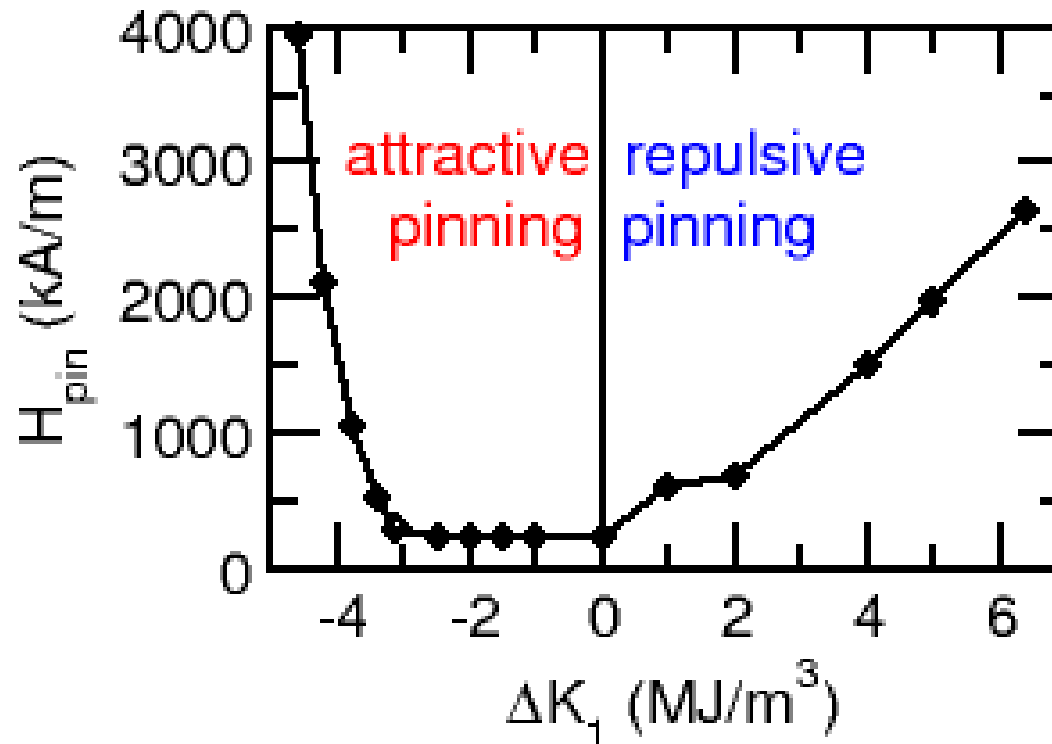
- Cells: $D = 125$ nm
- Intercellular phase:
Thickness: $t = 10$ nm
Anisotropy: $K_1 = 9$ MJ/m³



Domain wall depinning



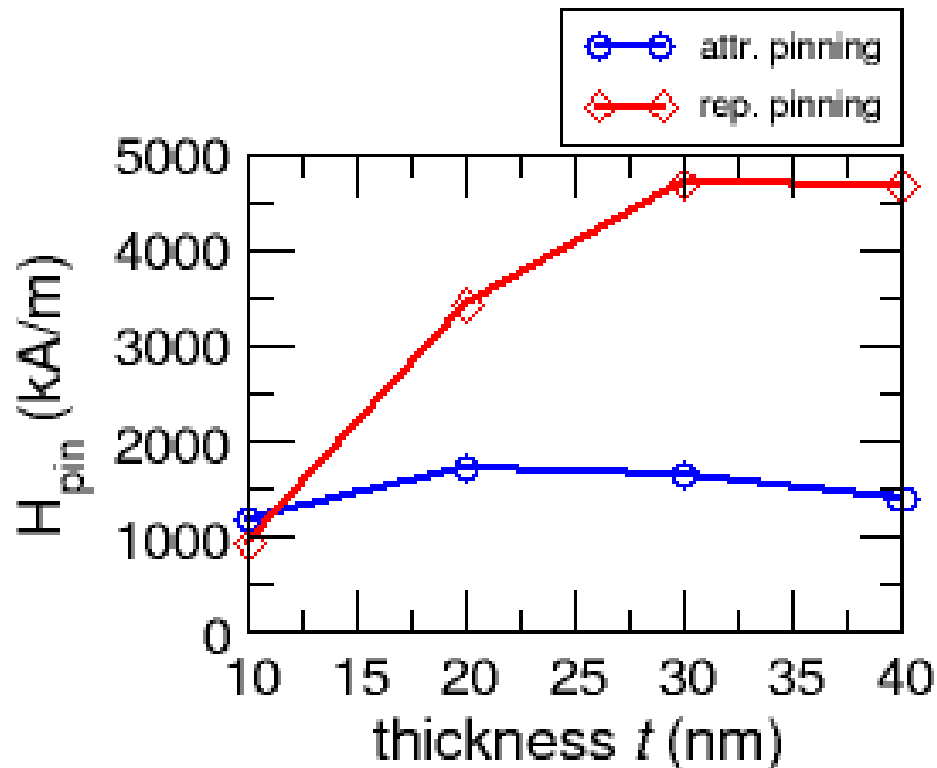
Pinning Fields



- Linear behaviour in the regime of repulsive pinning in agreement with a simple analytical 1D-model by Kronmüller (IEEE Trans. Magn. MAG-20 (1984) 1569):

$$H_c^{\max} = \alpha(K_1^{phase} - K_1^{cells}) / M_s^{cells}$$

Variation of the Phase Thickness



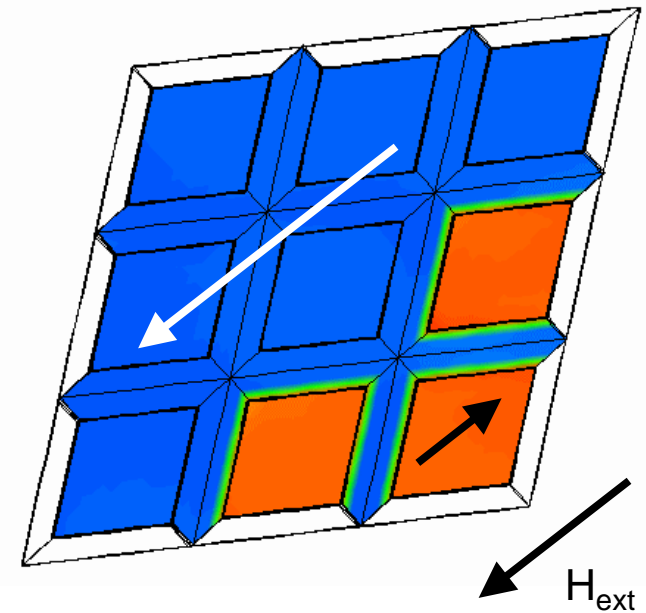
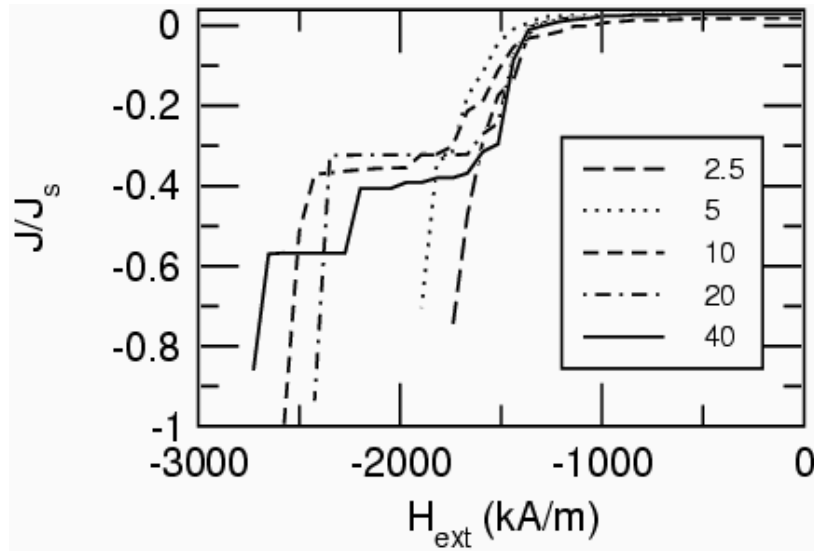
- Cells ($\text{Sm}_2\text{Co}_{17}$):
 $K_1 = 5 \text{ MJ/m}^3$

- Intercellular phase

Attractive pinning:
 $K_1 = 1.2 \text{ MJ/m}^3$

Repulsive pinning:
 $K_1 = 9.0 \text{ MJ/m}^3$

Thick Intercellular Phases



- Reversal of the whole intercellular phase
- Nucleation field of the cells determines H_c

Summary

- **Different pinning mechanisms depending on the composition**
- **Attractive Pinning:**
Pinning field decreases for increasing thickness of the intercellular phase
- **Repulsive pinning:**
linear dependence of the pinning field on anisotropy
maximum pinning field cannot be increased with increasing thickness of the intercell. phase

Acknowledgements

- **HITEMAG project**

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